

## OPERATIONAL STABILITY OF RUBIDIUM AND CESIUM FREQUENCY STANDARDS

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### INTRODUCTION

In the course of testing various rubidium and cesium frequency standards under operational conditions for use in NASA tracking stations, about 55 unit-years of relative frequency measurements for averaging times from 10 to  $10^7$ s have been accumulated at Goddard Space Flight Center (GSFC). Statistics on the behavior of rubidium and cesium standards under controlled laboratory conditions have been published by many institutions (see example, Ref. 1), but it was not known to what extent the lesser controlled environments of NASA tracking stations affected the performance of the standards. The purpose of this report is to present estimates of the frequency stability of rubidium and cesium frequency standards under operational conditions based on the data accumulated at GSFC.

Table 1.  
Atomic Frequency Standards Used  
in Experiments.

Serial no. or designation	Manufacturer
Rb 107	Varian Associates
Rb 136	Varian Associates
Rb 138	Varian Associates
Cs 110	Hewlett-Packard Co.
Cs 136	Hewlett-Packard Co.
Cs 137	Hewlett-Packard Co.
Cs 138	Hewlett-Packard Co.
Cs 139	Hewlett-Packard Co.
Cs 152	Hewlett-Packard Co.
Cs 182	Hewlett-Packard Co.
Cs 185	Hewlett-Packard Co.
Cs 186	Hewlett-Packard Co.
HM:	
H-10 no. 2	Varian Associates
NX-1	(a)

<sup>a</sup>An experimental hydrogen maser developed at GSFC. See Ref. 2.

## DATA DESCRIPTION

The three rubidium gas cells (designated Rb) and nine cesium beam frequency standards (C's) on which the measurements were made, as well as the two hydrogen masers (HM) used as references for many of the tests, are listed in Table 1 along with their serial numbers or designations and their manufacturers. During the tests the standards were kept in a laboratory at GSFC. Except for the shielding built into the standards themselves, there was no special control of the ambient magnetic, electric, vibration, and temperature conditions. The ambient magnetic and electric conditions were typically noisy. The standards were driven by ac power and were in no way isolated by transformers. Vibration from nearby air-conditioning equipment and from trucks at a nearby loading platform was not shielded in any way. The ambient temperature was typically between 298 and 303 K. There were, however, several brief excursions to temperatures as low as 291 K and as high as 313 K, due to equipment problems. These conditions are less controlled than those in the NASA tracking stations. Hence the stabilities of the standards when operating in the tracking stations should be at least as good as the stabilities calculated in this paper.

The measurements made on the standards consisted of average relative frequency measurements for varying averaging times. In some of the data sets, average relative frequency measurements were missing or were bad because of ac power failure or recorder failure. All such points were a posteriori linearly interpolated from the nearest earlier (in epoch time) good average relative frequency measurement and the nearest later (in epoch time) good average relative frequency measurement.

The total number of measurements made for all types of data used in this report is given in Table 2. Data sets are said to be of the same type when the following parameters are the same for each set: test unit,<sup>1</sup> reference unit, duration or averaging time  $\tau_0$  of each average relative frequency measurement, and dead time  $d$  between successive measurements (that is, the time during which no measurement was taken). The servo time constants are indicated only for the cesium standards and only when  $\tau_0 \leq 1000$  s. The difference in effect of a 10- and a 60-s time constant for  $\tau_0 \leq 3600$  s can be neglected because the time constants in such cases are too small with respect to standards to have an appreciable effect. The rubidium standards tested all have a fixed servo time constant which is on the order of 1 ms.

Neither temperature effects nor long-term frequency drift was removed from the data before analysis because the object of the tests was to measure the stability of the frequency standards under operational conditions, where both temperature fluctuations and long-term frequency drift are present.

<sup>1</sup>Although there are sometimes significant differences in the frequency stabilities of various rubidium standards, the three rubidium standards listed in Table 1 all had mutually close stabilities. For this reason, these rubidium standards will be considered to be identical. Because the nine cesium standards listed in Table 1 all had mutually close stabilities, they too will be considered to be identical.

Table 2  
Average Relative Frequency Data Sets.

Type of data				Number of data sets $m$	Number of measurements	
Test unit	Reference unit	Averaging time $\tau_0, s$	Dead time $d, s$		Total <sup>a</sup>	Interpolated
Rb	Rb	3 600	0.0	2	3 090	0
Rb	Cs (10-s TC)	10	2.3	1	1 076	18
Rb	Cs (10-s TC)	100	2.2	1	538	1
Rb	Cs (10-s TC)	1 000	2.7	1	223	0
Rb	HM	10	2.3	13	8 473	67
Rb	HM	100	2.2	10	6 405	16
Rb	HM	1 000	2.7	9	5 126	15
Rb	HM	3 600	.0	7	13 320	308
Cs	Cs	3 600	.0	3	8 851	263
Cs (10-s TC)	HM	10	.2	8	4 841	0
Cs (10-s TC)	HM	100	.2	8	4 871	11
Cs (10-s TC)	HM	1 000	.2	8	4 787	25
Cs (60-s TC)	HM	10	.2	8	4 634	0
Cs (60-s TC)	HM	10	2.3	3	1 904	2
Cs (60-s TC)	HM	100	.2	8	4 706	0
Cs (60-s TC)	HM	100	2.2	3	2 496	0
Cs (60-s TC)	HM	1 000	.2	8	4 804	3
Cs (60-s TC)	HM	1 000	2.7	1	692	0
Cs	HM	3 600	.0	13	37 404	1391
Cs	HM	604 800	.0	1	88	6

TC = time constant.

<sup>a</sup>Total number of measurements for all  $m$  data sets, including the interpolated measurements.

## STATISTICAL ANALYSIS

Let there be given a set of  $m$  identical test frequency standards and a set of  $m$  identical reference frequency standards. Let  $\phi_n(t)$ ,  $1 \leq n \leq m$ , denote the instantaneous fluctuations (measured in time units) of the epoch time output of the  $n$ th reference standard. Let  $y_n(t)$  be the instantaneous (fractional) frequency fluctuation of the  $n$ th test standard compared with the  $n$ th reference standard; i.e.,

$$y_n(t) \equiv \frac{d\phi_n(t)}{dt} \quad (1)$$

Let  $\bar{y}_n(t)$  be the average relative (fractional) frequency fluctuation of the  $n$ th test standard compared with the  $n$ th reference standard:

$$\bar{y}_n(t) = \frac{1}{\tau} \int_t^{t+\tau} y_n(t) dt = \frac{\phi_n(t+\tau) - \phi_n(t)}{\tau} \quad (2)$$

The constant  $\tau$  is called the averaging time of  $y(t)$ . The Allan standard deviation  $\sigma(2, T, \tau)$  of the frequency fluctuations of the set of test standards compared with the set of reference standards is defined to be (Ref. 3)

$$\sigma(2, T, \tau) = \sqrt{\frac{1}{m} \sum_{n=1}^m \langle \text{var} [\bar{y}_n(t+T) - \bar{y}_n(t)] \rangle} \quad (3)$$

where the symbol  $\langle \rangle$  denotes infinite epoch time average. The analysis of all data listed in Table 2 consisted in the calculation of an estimate, which is denoted by  $s(2, T, \tau)$  in the following manner.

Taking any type of data from Table 2, let the number of average relative frequency measurements in the  $n$ th data set,  $1 \leq n \leq m$ , be  $m_n$ . Denote this  $n$ th set of average relative frequency measurements by  $y_n(i)$ ,  $i=1, \dots, m_n$ . For  $i=1, 2, \dots, m_n-1$ , denote the variance of the two average relative frequency measurements  $y_n(i)$  and  $y_n(i+1)$

$$v_n(i) = \frac{[\bar{y}_n(i+1) - \bar{y}_n(i)]^2}{2} \quad (4)$$

The square root of the average over both  $i$  ( $1 \leq i \leq m_n - 1$ ) and  $n$  ( $1 \leq n \leq m$ ) of these  $v_n(i)$  is the desired estimate of  $\sigma(2, \tau_0 + d, \tau_0)$ :

$$s(2, \tau_0 + d, \tau_0) = \sqrt{\frac{\sum_{n=1}^m \sum_{i=1}^{m_n-1} v_n(i)}{\sum_{n=1}^m (m_n - 1)}} \quad (5)$$

From the original data sets  $y_n$   $\sum_{i=1}^{m_n}$ ,  $1 \leq n \leq m$ , new data sets with averaging time  $\tau_1 = 2\tau_0$  and dead time  $d$  (assumed small with respect to  $\tau_0$ ) can be approximated by defining

$$\bar{y}_n(i; 1) = \frac{\bar{y}_n(i+1) + \bar{y}_n(i)}{2} \quad (6)$$

$i = 1, 2, \dots, m_n - 1$  and  $n = 1, 2, \dots, m$ . Denote the variance of  $y_n(i; 1)$  and  $y_n(i+2; 1)$  by  $v_n(i; 1)$ :

$$v_n(i; 1) = \frac{[\bar{y}_n(i+2; 1) - \bar{y}_n(i; 1)]^2}{2} \quad (7)$$

$i = 1, 2, \dots, m_n - 3$  and  $n = 1, 2, \dots, m$ . Estimate  $\sigma(2, \tau + d, \tau_1)$  by<sup>2</sup>

$$s(2, \tau_1 + d, \tau_1) = \sqrt{\frac{\sum_{n=1}^m \sum_{i=1}^{m_n-3} v_n(i; 1)}{\sum_{n=1}^m (m_n - 3)}} \quad (8)$$

Let  $k$  be the exponent of the largest power of 2 contained in any of the  $m_n$ ,  $1 \leq n \leq m$ . For  $j = 2, 3, \dots, k-1$ , the data set  $\{\bar{y}_n(i; j)\}_{i=1}^{m_n-2^j+1}$  with averaging time  $\tau_j = 2^j \tau_0$  and dead time  $d$  is successively calculated from the data set  $\{\bar{y}_n(i; j-1)\}_{i=1}^{m_n-2^{j-1}+1}$  by pairwise averaging:

$$\bar{y}_n(i; j) = \frac{\bar{y}_n(i+2^{j-1}; j-1) + \bar{y}_n(i; j-1)}{2} \quad (9)$$

$i = 1, 2, \dots, m_n - 2^j + 1$ ;  $n = 1, 2, \dots, m$ ;  $j$  fixed. Denote the variance of  $\bar{y}_n(i; j)$  and  $\bar{y}_n(i+2^j; j)$  by  $v_n(i; j)$ :

$$v_n(i; j) = \frac{[\bar{y}_n(i+2^j; j) - \bar{y}_n(i; j)]^2}{2} \quad (10)$$

<sup>2</sup>Throughout this paper the convention is adopted that whenever a summand, e.g.,  $m_n - 3$  in  $\sum_{n=1}^m (m_n - 3)$ , is less than zero, it is treated as zero; and whenever a summation, e.g.,  $\sum_{i=1}^{m_n-3} v_n(i; 1)$ , has an upper limit that is less than the lower limit, it also is treated as zero.

$i = 1, 2, \dots, m_n - 2^{j+1} + 1$  and  $n = 1, 2, \dots, m$ . Estimate  $\sigma(2, \tau_j + d, \tau_j)$  by

$$s(2, \tau_j + d, \tau_j) = \sqrt{\frac{\sum_{n=1}^m \sum_{i=1}^{m_n - 2^{j+1} + 1} v_n(i; j)}{\sum_{n=1}^m (m_n - 2^{j+1} + 1)}} \quad (11)$$

An example of this procedure for zero dead time is presented in Figure 1. The quantity  $v$  represents the variance between the ordinates of the two lines to which the dotted line near  $v$  points.

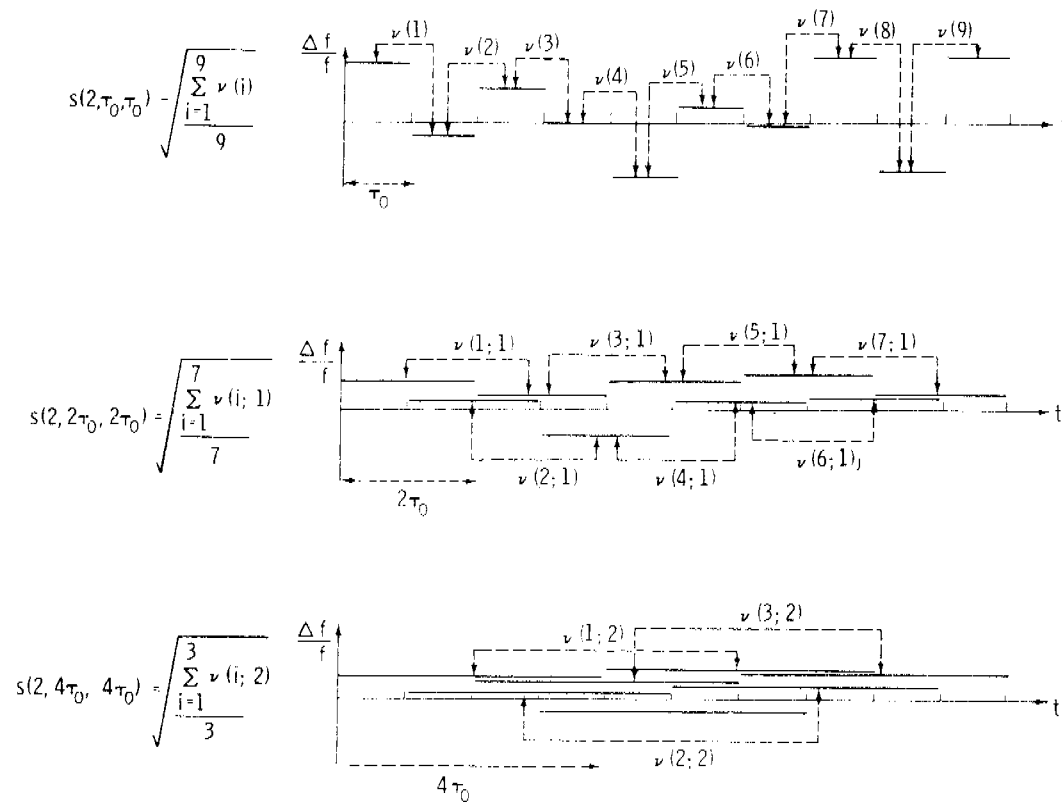


Figure 1. Calculation of  $s(2, \tau, \tau)$ .

## RESULTS

For each type of data listed in Table 2 and for each averaging time  $\tau_j = 2^j \tau_0$ ,  $0 \leq j \leq k-1$  ( $\tau_0$  and  $k$  change with the type of data), the estimate  $s(2, \tau_j)$  of  $\sigma(2, \tau_j + d, \tau_j)$  was calculated.<sup>3</sup> The results are presented in Table 3 and Figure 2 for all data involving a rubidium standard as either the test or the reference unit and in Table 4 and Figure 3 for the cesium versus cesium and cesium versus hydrogen maser data.

In order to use the data in Tables 3 and 4 to estimate the frequency stability of the rubidium and cesium standards tested, rather than the relative frequency stability of a comparison of two of these standards or of a comparison of one of these standards to a hydrogen maser, the following procedure is used. Denote the Allan standard deviations of the test standard versus a hypothetical perfect standard, the reference standard by  $\sigma_T(2, \tau + d, \tau)$ ,  $\sigma_R(2, \tau + d, \tau)$ , and  $\sigma_{T-R}(2, \tau + d, \tau)$  respectively. Because the variances  $\sigma_T^2(2, \tau + d, \tau)$  and  $\sigma_R^2(2, \tau + d, \tau)$  are linear functions (in fact, weighted integrals) of the respective power spectral densities of the test and reference standards (Ref. 3); and because the power spectral density of the comparison of two frequency standards is the sum of the power spectral densities of each of the standards, the following relation occurs:

$$\sigma_{T-R}^2(2, \tau + d, \tau) = \sigma_T^2(2, \tau + d, \tau) + \sigma_R^2(2, \tau + d, \tau) \quad (12)$$

For comparisons of two identical standards (rubidium standard versus rubidium standard and cesium standard versus cesium standard),  $\sigma_R(2, \tau + d, \tau) = \sigma_T(2, \tau + d, \tau)$ . Hence, from relation (12),

$$\sigma_T(2, \tau + d, \tau) = \frac{\sigma_{T-R}(2, \tau + d, \tau)}{\sqrt{2}} \quad (13)$$

For all data for which a hydrogen maser was used as a reference, it is assumed that the instabilities of the maser were sufficiently small so as to have

$$\sigma_T(2, \tau + d, \tau) \approx \sigma_{T-R}(2, \tau + d, \tau) \quad (14)$$

The normalized standard deviation  $\sigma_T(2, \tau, \tau)$  can be calculated from  $\sigma_T(2, \tau + d, \tau)$  by the relation

$$\sigma_T(2, \tau, \tau) = \frac{\sigma_T(2, \tau + d, \tau)}{\sqrt{B_2(r, \mu)}} \quad (15)$$

where  $B_2(\tau, \mu)$  is a bias function (defined in ref. 4);  $r = (\tau + d)/\tau$ ; and  $\mu$ , representing the type of noise of the standard for the fixed averaging time  $\tau$  and fixed dead time  $d$ , is determined from

$$\sigma_T(2, \tau + d, \tau) \propto \tau^{\mu/2} \quad (16)$$

<sup>3</sup>The analysis was carried out by programs E00016 and E00036 of the GSFC Computer Program Library. Program E00016 is for input relative phase data; program E00036 is for input relative frequency data. Although program E00016 reads relative phase data as input, its output is the Allan Standard deviation of relative frequency  $s(2, \tau_j + d, \tau_j)$  defined in eqs. (5), (8), and (11). These two programs are based on a program written by David W. Allan of the National Bureau of Standards, Boulder, Colo.





Table 4.  
Cesium Standard Frequency Stability.

Type of data				Type of data				Type of data					
Test unit	Reference unit	$\tau, s$	$d, s$	Test unit	Reference unit	$\tau, s$	$d, s$	Test unit	Reference unit	$\tau, s$	$d, s$	$s(2, \tau + d, \tau), \times 10^{-12}$	$s(2, \tau + d, \tau), \times 10^{-12}$
Cs	Cs	3 600	0.0	Cs (60-s TC)	HM	256 000	0.2	Cs (60-s TC)	HM	1 000	0.2	1.825	2.199
		7 200	.0			10	.2			2 000	.2	5.391	1.523
		14 400	.0			20	.2			4 000	.2	4.524	1.099
		28 800	.0			40	.2			8 000	.2	4.510	.790
		57 600	.0			80	.2			16 000	.2	4.529	.590
		115 200	.0			160	.2			32 000	.2	4.149	.435
		230 400	.0			320	.2			64 000	.2	3.149	.360
		460 800	.0			640	.2			128 000	.2	2.287	.343
		921 600	.0			1 280	.2			256 000	.2	1.786	.375
		1 843 200	.0			2 560	.2			512 000	.2	1.201	3.885
		3 686 400	.0			5 120	.2			1 024 000	.2	6.080	2.963
		7 372 800	.0			2 048	.2			409 600	.2	5.321	2.170
		14 745 600	.0			819 200	.2			16 384 000	.2	5.875	1.542
29 491 200	.0	32 784 000	.2	65 539 200	.2	5.912	1.042						
58 982 400	.0	131 168 000	.2	262 176 000	.2	5.882	.672						
117 964 800	.0	524 704 000	.2	1 096 704 000	.2	5.875	.409						
235 929 600	.0	2 149 408 000	.2	8 787 200 000	.2	5.912	.262						
471 859 200	.0	4 298 816 000	.2	17 574 400 000	.2	5.039	.732						
943 718 400	.0	8 597 632 000	.2	35 148 800 000	.2	3.636	.615						
1 887 436 800	.0	17 195 264 000	.2	70 297 600 000	.2	2.217	.572						
3 774 873 600	.0	34 390 528 000	.2	140 595 200 000	.2	1.943	.566						
7 549 747 200	.0	68 781 056 000	.2	281 190 400 000	.2	4.149	.570						
15 099 494 400	.0	137 562 112 000	.2	562 380 800 000	.2	3.687	.556						
30 198 988 800	.0	275 124 224 000	.2	1 124 761 600 000	.2	2.962	.555						
60 397 977 600	.0	550 248 448 000	.2	2 249 523 200 000	.2	2.144	.590						
120 795 955 200	.0	1 100 496 896 000	.2	4 499 046 400 000	.2	1.535	.612						
241 591 910 400	.0	2 200 993 792 000	.2	8 998 092 800 000	.2	1.072	.464						
483 183 820 800	.0	4 401 987 584 000	.2	17 996 185 600 000	.2	.760	.337						
966 367 641 600	.0	8 803 975 168 000	.2	35 992 371 200 000	.2	.580	.239						
1 932 735 283 200	.0	17 607 950 336 000	.2	71 984 742 400 000	.2	.595	.181						
3 865 470 566 400	.0	35 215 900 672 000	.2	143 969 484 800 000	.2	.599	.167						
7 730 941 132 800	.0	70 431 801 344 000	.2	287 938 969 600 000	.2	5.699	.114						
15 461 882 265 600	.0	140 863 602 688 000	.2	575 877 939 200 000	.2	5.535							
30 923 764 531 200	.0	281 727 205 376 000	.2	1 151 754 878 400 000	.2	4.840							
61 847 529 062 400	.0	563 454 410 752 000	.2	2 303 509 756 800 000	.2	3.690							
123 695 058 124 800	.0	1 126 908 821 504 000	.2	4 607 019 513 600 000	.2	2.750							
247 390 116 249 600	.0	2 253 817 643 008 000	.2	9 214 039 027 200 000	.2	1.931							
494 780 232 499 200	.0	4 507 635 286 016 000	.2	18 428 078 054 400 000	.2	1.318							
989 560 464 998 400	.0	9 015 270 572 032 000	.2	36 856 156 108 800 000	.2	1.003							
1 979 120 929 996 800	.0	18 030 541 144 064 000	.2	73 712 312 217 600 000	.2	.883							
3 958 241 859 993 600	.0	36 061 082 288 128 000	.2	147 424 624 435 200 000	.2								

It is of interest to note that for the  $\tau$  and  $\mu$  of the data analyzed in this report,  $B_2(\tau, \mu)$  differs from unity by less than 0.1 percent and can be ignored. Hence, for the data in this report,

$$\sigma_T(2, \tau, \tau) \approx \sigma_T(2, \tau + d, \tau) \quad (17)$$

Of course, relation (17) is an exact equality whenever  $d = 0$ .

Using the estimates  $s(2, \tau + d, \tau)$  of  $\sigma_{T-R}(2, \tau + d, \tau)$  from Figures 2 and 3 in relations (13) and (14) and using relation (17), the standard deviations  $\sigma_T(2, \tau, \tau)$  of the rubidium and cesium standards tested can be estimated. These estimates of  $\sigma_T(2, \tau, \tau)$  are presented in Figure 4 as the "operational environment" curves. Also shown in Figure 4 are curves taken from References 1 and 5 representing the performance of rubidium and cesium standards in a "controlled environment." By "controlled environment" is meant an experimental environment shielded from magnetic, electric, vibration, and temperature effects much more than the "operational" environment in which the data presented in Figures 2 and 3 were taken.<sup>4</sup> The upper curve for rubidium standards under a controlled environment in Figure 4 is taken from Reference 5 and represents the measured performance of Varian rubidium standards under controlled conditions. The lower curve for rubidium standards under a controlled environment and the curve for cesium standards under a controlled environment in Figure 4 are taken from Reference 1 and represent the measured performance of Hewlett-Packard rubidium and cesium standards under controlled conditions.

## CONCLUSIONS

From Figure 4 it is apparent that an operational environment degrades the performance of the rubidium standards (by up to one order of magnitude) for frequency averaging times between 10 and  $10^3$  s and that it degrades the performance of the cesium standards (by up to one order of magnitude) for frequency averaging times between  $3 \times 10^4$  and  $2 \times 10^7$  s. For all other averaging times in the range covered by the data in Figure 4, the stabilities of the standards are not degraded by the operational conditions.

## ACKNOWLEDGEMENTS

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<sup>4</sup>For Reference 1, these conditions were verified by G. M. R. Winkler (private communication). It is assumed that the data presented in Reference 5 were obtained in a similarly controlled environment.

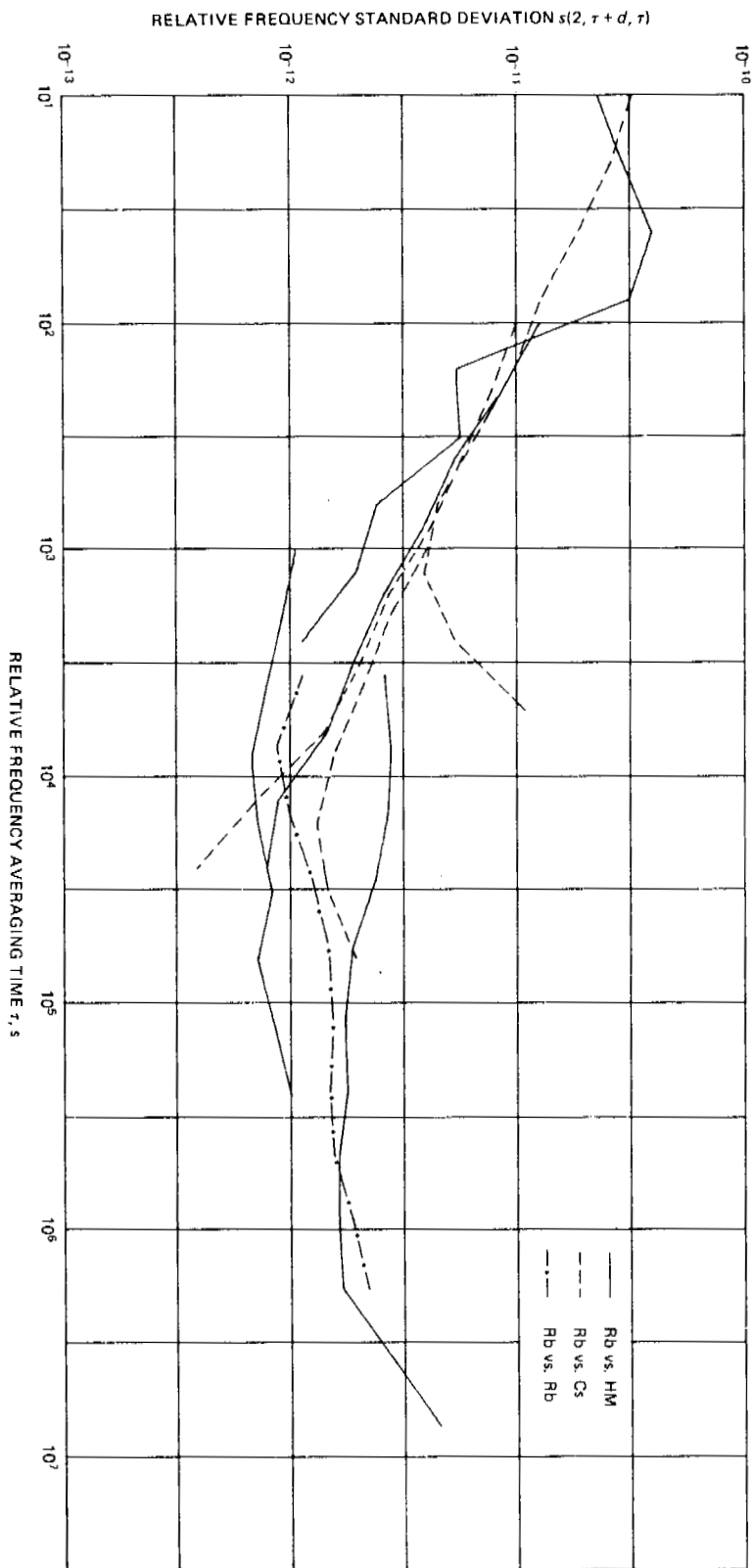


Figure 2. Rubidium standard relative frequency stability.

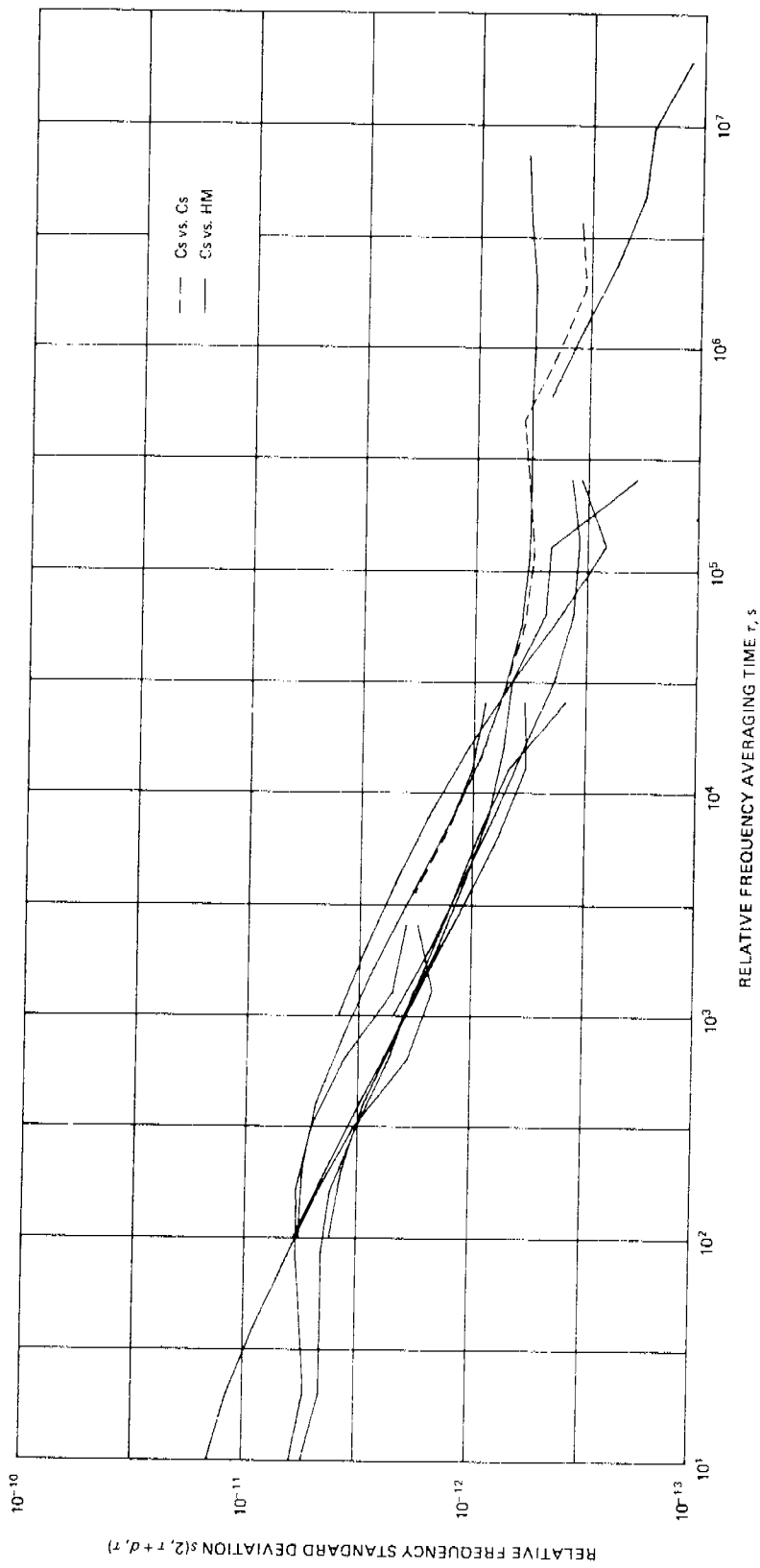


Figure 3. Cesium standard relative frequency stability.

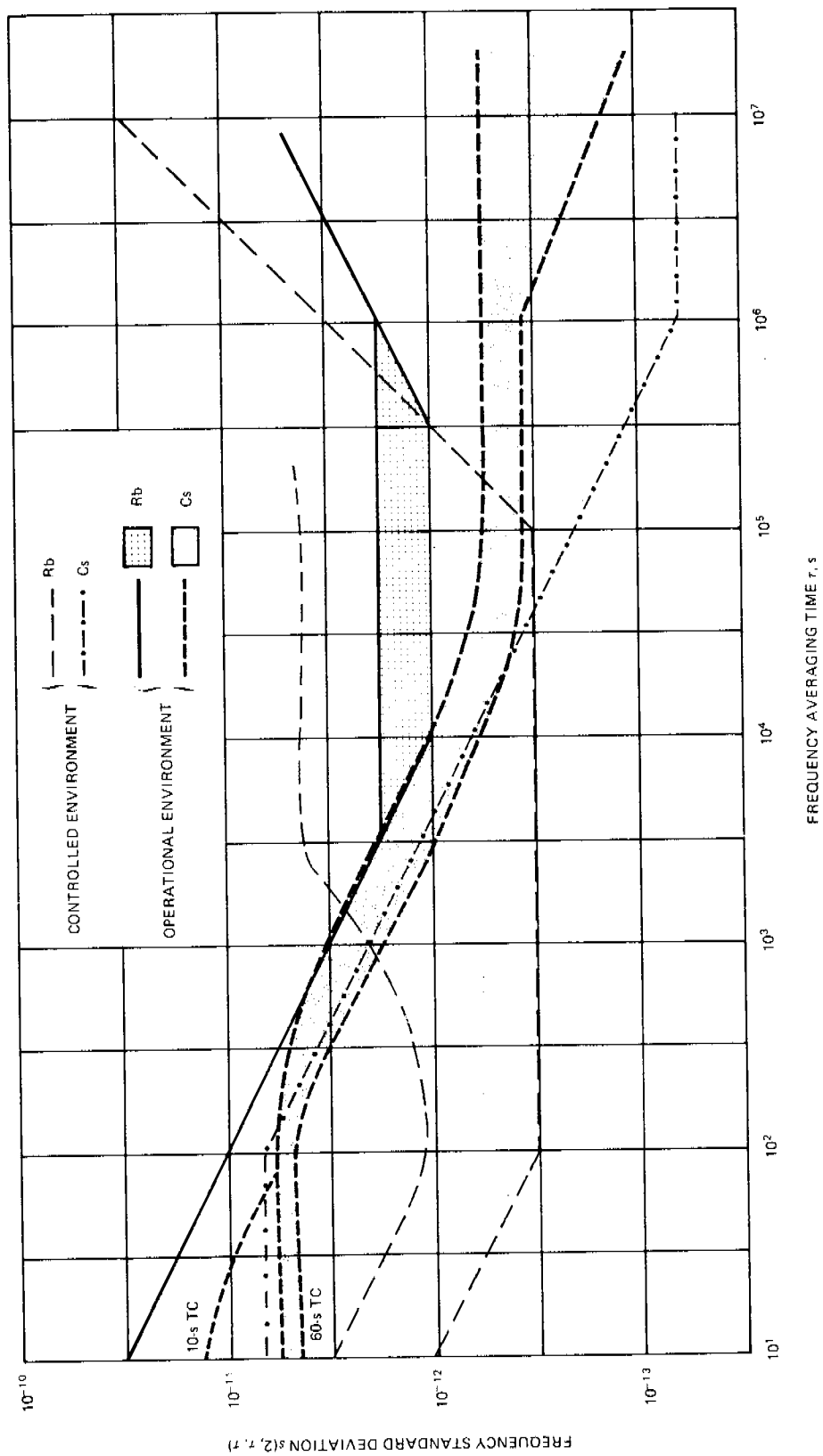


Figure 4. Rubidium and cesium standard frequency stabilities.

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